

EFFECTS OF PEENING ON MAGNETIC PROPERTIES IN SOFT MAGNETIC ALLOYS

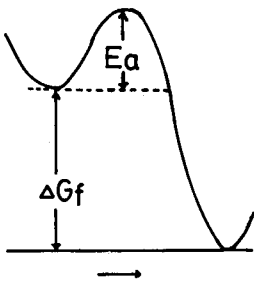
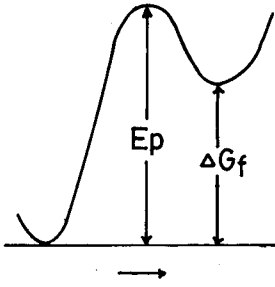
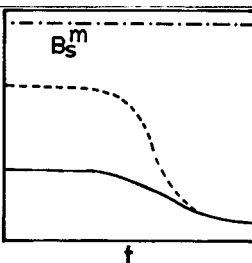
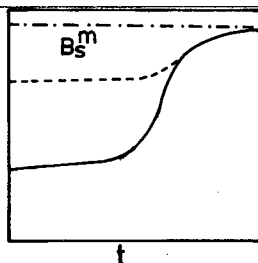
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Introduction

Various properties, such as magnetic, electric and mechanical properties, have been reported in liquid-quenched metallic glasses (1 - 6). However, it is difficult to control the properties, homogeneously, because cooling condition radically affects the properties. Aging, close to equilibrium, is a good way of homogenizing magnetic properties of metallic glasses (7 - 8). This homogenization relaxes heterogeneous elastic stress induced by liquid-quenching, although the aging decays ductility (9) (see Table 1).

Peening is effective for increasing the compressive stress, homogeneously. This stress enlarges resistances to fatigue (10) and stress-corrosion (11). Furthermore, the peening may not only enhance the compressive stress homogeneously on the surface, but also disorder the structure of metallic glasses, homogeneously (Table 1).

Table 1: Rate process of peening and aging

	A g i n g	P e e n i n g
S t r u c t u r e C h a n g e	H o m o g e n i z a t i o n	
	① Stabilization ② Ordering ③ Relaxation	① unstabilization ② disordering ③ random distribution
R a t e P r o c e s s		
D r i v i n g F o r c e	thermal energy	collision energy
B s C h a n g e		

The dependence of the saturated magnetic properties of metallic glasses on the cooling rate has been investigated (12). It was seen that the fast-quenched and unrelaxed metallic glasses exhibit high magnetic flux density B_s . With this experience, we have undertaken the present study to observe the effect of peening on the magnetic properties in soft magnetic alloys

Experimental procedure

The foil samples of Fe-15at%Ni-10at%Si-15at%B were prepared by liquid-quenching with a twin-type piston-anvil apparatus under a protective Ar-5% H_2 atmosphere (13 - 14). The samples were quenched from approximately 1700 K. The speed of the piston was about 0.12 m/s. Cooling rate, indicated by a parameter, D , was changed by controlling the thickness of the samples. Unrelaxed and relaxed glass samples are prepared by fast- and slow-cooled liquid-quenching below a critical thickness to which the glass can be prepared.

The peening is performed by an apparatus shown in Fig. 1. The nozzle diameter was 8.0 mm. The velocity of the air at the nozzle was 195.5 m/s. The distance between the nozzle and the specimen was 20 mm. The peening angle was 30° to the specimen. The steel balls were made of SuJ 2 steel (HRC = 64); their mean diameter and weight were 0.4 ± 0.1 mm and 0.68 mg, respectively, and they were supplied at a rate of $2.660 \text{ mm}^{-1} \text{ s}^{-1}$.

The structure of the the samples was examined by means of X-ray diffraction. B_s was measured with a B-H curve tracer (BHV-3.0, Riken Denshi, Tokyo). We found that the magnetic flux density at 5 kOe was very slightly (1 to 5%) larger than that at 4 kOe. Thus, the value of the flux density at 5 kOe was

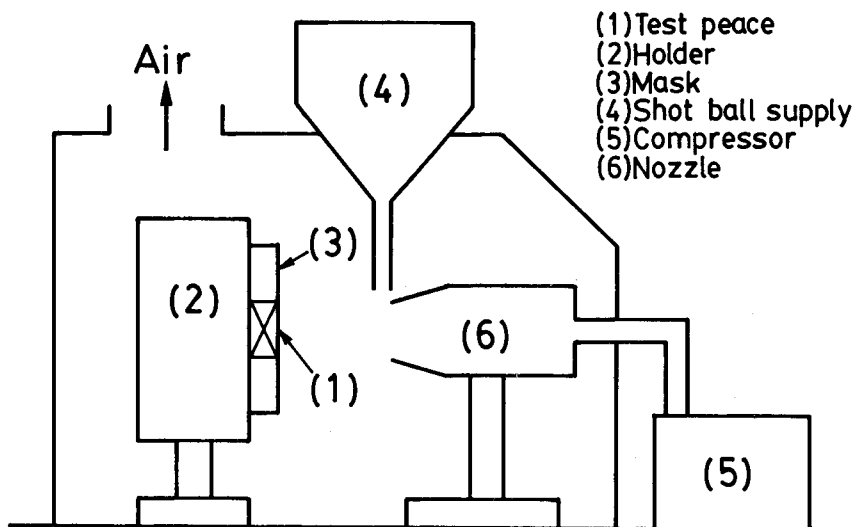


Fig. 1 Schematic peening apparatus.

adopted for B_s . Since it was difficult to obtain the corrected value of the coercive force H_c , the reduced coercive force H_c^r [$(H_c - H_c^0)/H_c^0$] was taken as an indicator of coercive force. Here, H_c^0 is the H_c before peening.

Results and discussions

Besides the glass, the changes in the magnetic properties are studied for Fe-Si crystal alloy. This alloy is commercially used as a typical soft magnetic steel. Figure 2 shows B_s and H_c^r against peening time t of Fe-3%Si crystal alloy. B_s and H_c^r increase with an increase in peening time. It is obvious that the increase in peening time t enhances the B_s and H_c^r . Cold-rolling increases B_s and H_c^r , too (Fig. 3). The cold-working is concluded to increase B_s and H_c^r . The disordering increases B_s and H_c^r . Thus, the peening deduces to increase B_s and H_c^r .

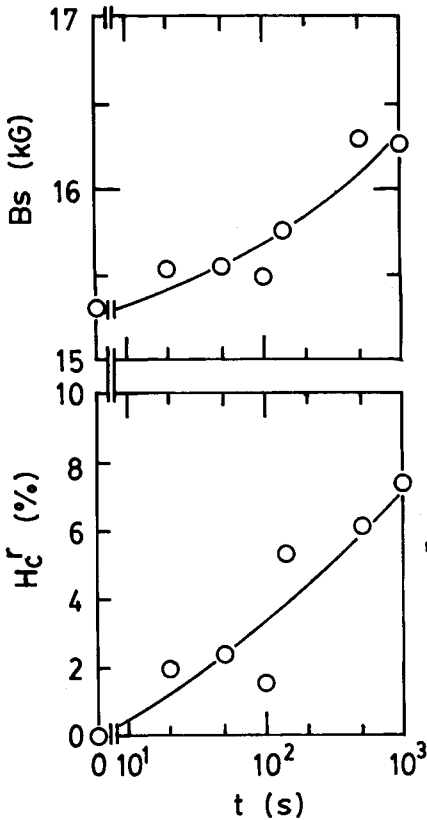


Fig. 2: Change in saturated magnetic flux density B_s and reduced coercive force H_c^r with peening time of Fe-3%Si crystal alloy.

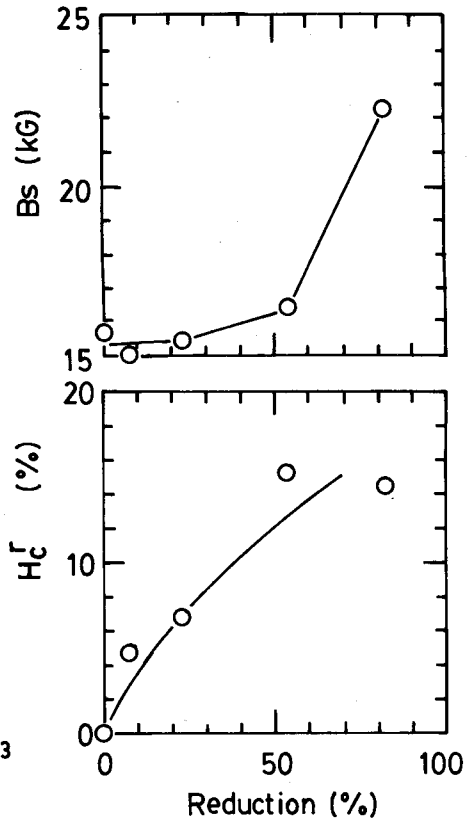


Fig. 3: Change in saturated magnetic flux density B_s and reduced coercive force H_c^r with reduction of cold-rolled Fe-3%Si crystal alloy.

Figure 4 shows B_s against peening time t of liquid-quenched Fe-15at%Ni-10at%Si-15at%B alloy samples. Before peening, at $T = 0$ sec, the thinner the specimens (the faster the cooling rate), the larger the B_s becomes (12). Compared with the crystal sample, the glassy samples show high B_s . For the fast and slow cooled samples, B_s increases with an increase in peening time. It is obvious that the increase in peening time t enhances the B_s . The largest value of B_s is 9.08 kG ($t = 1000$ s, $D = 0.105$ mm) of the fast cooled glass.

A relaxed and clustered metallic glass has prismatic structure (15 - 16). Since electrons of metalloids atoms move to holes of transition metal atoms in the relaxed state, the magnetic moments of the transition metal-metalloid glasses are in general lower than those of alloys without metalloids. Since the relaxed glass is obtained by slow quenching, B_s of the fast cooled specimen is higher than that of the slow cooled specimen (see B_s at $t = 0$ sec in Fig. 4), i.e. the relaxed glass has the lower B_s . Thus, the peening dependence of B_s in the glassy state in Fig. 4 is understood.

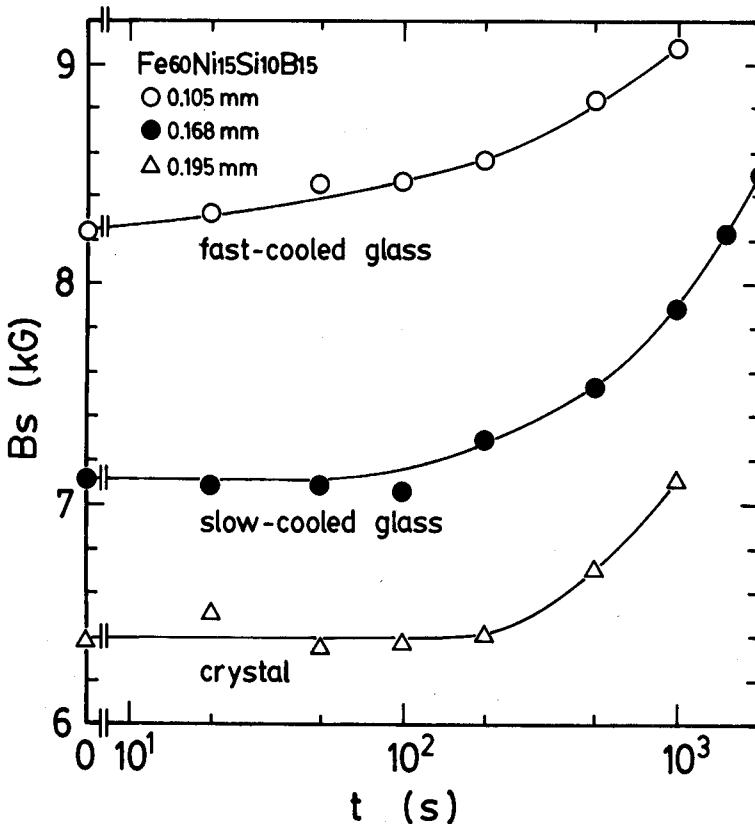


Fig. 4 Change in saturated magnetic flux density B_s with peening time of liquid-quenched Fe-15at%Ni-10at%Si-15at%B alloy.

Figure 5 shows the reduced coercive force H_c^r [$(H_c - H_c^0)/H_c^0$] with the peening time of the liquid-quenched Fe-15at%Ni-10at%Si-15at%B alloy. H_c^r slightly increases for the crystal alloy, which is prepared by slow-liquid-quenching. For the glasses, H_c^r increases with an increase in the peening time. It is obvious that the increase in peening time t enhances the H_c^r . On the other hand, the cooling condition does not affect the value of the H_c^r for the glasses.

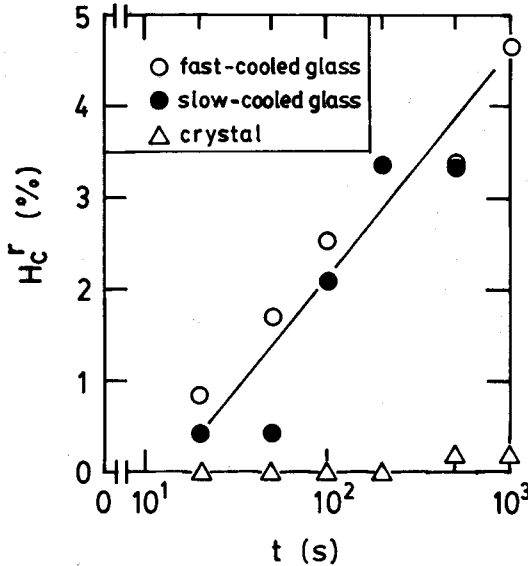


Fig. 5 Change in reduced coercive force H_c^r [$(H_c - H_c^0)/H_c^0$] with the peening time of a liquid-quenched Fe-15at%Ni-10at%Si-15at%B alloy. Here, H_c^0 is the H_c before peening.

Rate process of peened glass

If a driving force of the peening is collision energy, a rate process is applied (see Table 1). Based on the rate process (17), the B_s change (X) is assumed to be expressed by a following equation in relation to the peening time (t ; sec).

$$X = 1 - \exp(-k t)^n \quad (1)$$

Here, k and n are constant. X is assumed to express

$$X = [(B_s - B_s^0)/(B_s^m - B_s^0)], \quad (2)$$

where B_s^m and B_s^0 are B_s of the peened glass at extremely long period of time and of as-quenched glass (or minimum value of B_s), respectively. Peening for both fast- and slow-cooled glasses, B_s approaches B_s^m (see Fig. 6). B_s^m of eq. (2) is 9.510 kG, when the correlation coefficient (F) of eq. (1) is maximum (F

= 0.9912 for the slow-cooled glass and $F = 0.9853$ for the fast-cooled glass) as shown in Fig. 7.

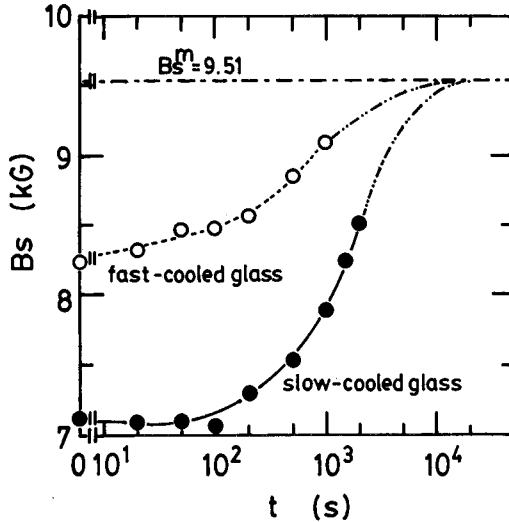


Fig. 6: Change in saturated magnetic flux density B_s with peening time t of liquid-quenched Fe-15at%Ni-10at%Si-15at%B alloy glasses.

From these results, X is expressed by a following equation of fast- and slow-cooled Fe-Ni-Si-B alloy glasses (see Fig. 8).

$$\log_{10} [-\ln(1-X)] = 0.650 \log t - 1.94 \quad (\text{fast-cooled glass}) \quad (3)$$

$$\log_{10} [-\ln(1-X)] = 0.984 \log t - 3.32 \quad (\text{slow-cooled glass}) \quad (4)$$

This linear plots confirms the assumption of eq. (1). B_s is plotted of the solid line in Fig. 8. B_s of the peened glass approaches B_s^m in Fig. 6.

Conclusion

The B_s and H_c changes with peening is studied for soft magnetic alloys. The peened alloys show the large B_s and the large H_c . If a driving force of the peening is collision energy, a rate process is applied for metallic glass. The parameter of X , $[(B_s - B_s^0)/(B_s^m - B_s^0)]$, is assumed to express, where B_s^m and B_s^0 are B_s of the peened glass at extremely long period of time and of as-quenched glass (or minimum value of B_s), respectively. Based on the rate process, X is expressed by a following equation of fast- and slow-cooled Fe-Ni-Si-B alloy glasses.

$$\log_{10} [-\ln(1-X)] = 0.650 \log t - 1.94 \quad (\text{fast-cooled glass})$$

$$\log_{10} [-\ln(1-X)] = 0.984 \log t - 3.32 \quad (\text{slow-cooled glass})$$

This linear plots confirms the assumption of the rate process. Peening for both fast- and slow-cooled glasses, B_s approaches B_s^m . B_s^m is 9.510 kG, when the correlation coefficient is maximum.

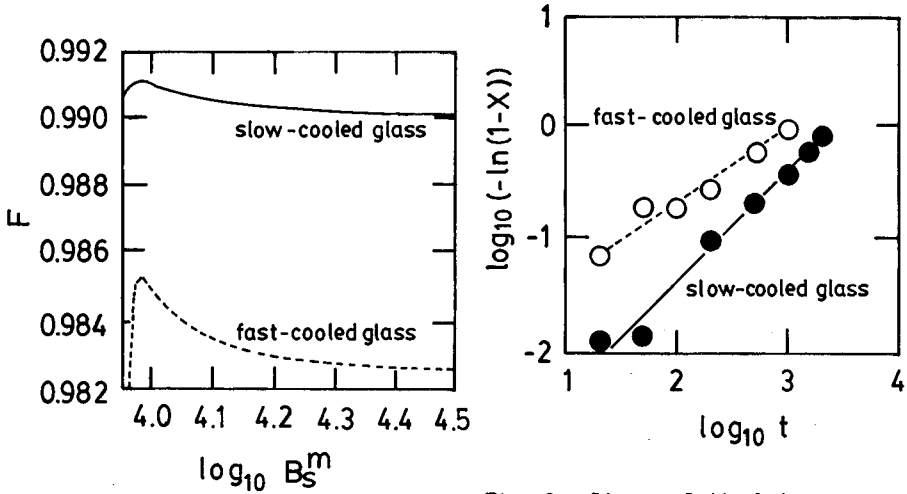


Fig. 7: Change in correlation coefficient (F) with B_s for fast- and slow-cooled glasses. B_s^m , which is the B_s at extremely fast-cooling, is obtained at the maximum value of F .

Fig. 8: Linear plots between $\log_{10}[-\ln(1-X)]$ and $\log t$. $X = [(B_s - B_s^0)/(B_s^m - B_s^0)]$, where B_s^m and B_s^0 are B_s of the peened glasses of extremely long period of time and as-quenched glass (or minimum value of B_s), respectively.

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